

# Ising chain: Thermal conductivity and first-principle validation of Fourier's law



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## Introduction

- Fourier's law [1] describes the heat diffusion rate through a macroscopic material in the direction of the flow ( $\mathbf{J} \propto -\nabla T$ );
- In this work two types of anisotropic planar rotators are studied through molecular dynamics;
- We approach, for a linear chain, the Ising limit via two different types of extremely anisotropic XY models (local and in the coupling), which allowed us to evaluate the validity of the Fourier's law;
- Furthermore, we better characterized the conductivity change for a more extended range of temperatures, resulting in the  $q$ -stretched exponential [2] instead of the  $q$ -Gaussian distribution [3].

## Models

- The Hamiltonians of the local and anisotropic coupling models are, respectively, given by

$$\mathcal{H}_{XY}^l = \sum_{i=1}^L \frac{p_i^2}{2} + \frac{1}{2} \sum_{\langle i,j \rangle} [1 - \cos(\theta_i - \theta_j)] + \epsilon_l \sum_{i=1}^L \sin^2 \theta_i \quad (1)$$

and

$$\mathcal{H}_{XY}^a = \sum_{i=1}^L \frac{p_i^2}{2} + \frac{1}{2} \sum_{\langle i,j \rangle} [1 + \epsilon_a \cos(\theta_i - \theta_j) - \epsilon_a \cos(\theta_i + \theta_j)]. \quad (2)$$

where  $\epsilon_l \in [0, \infty)$  and  $\epsilon_a \in [-1, 1]$ ;

- $\epsilon_a = \pm 1$  correspond to the Ising model along the  $y$  and  $x$  axes respectively, whereas  $\epsilon_a = 0$  recovers the standard isotropic XY-model (see Fig. 1);
- We have considered unit momenta of inertia and unit first-neighbor coupling constant;

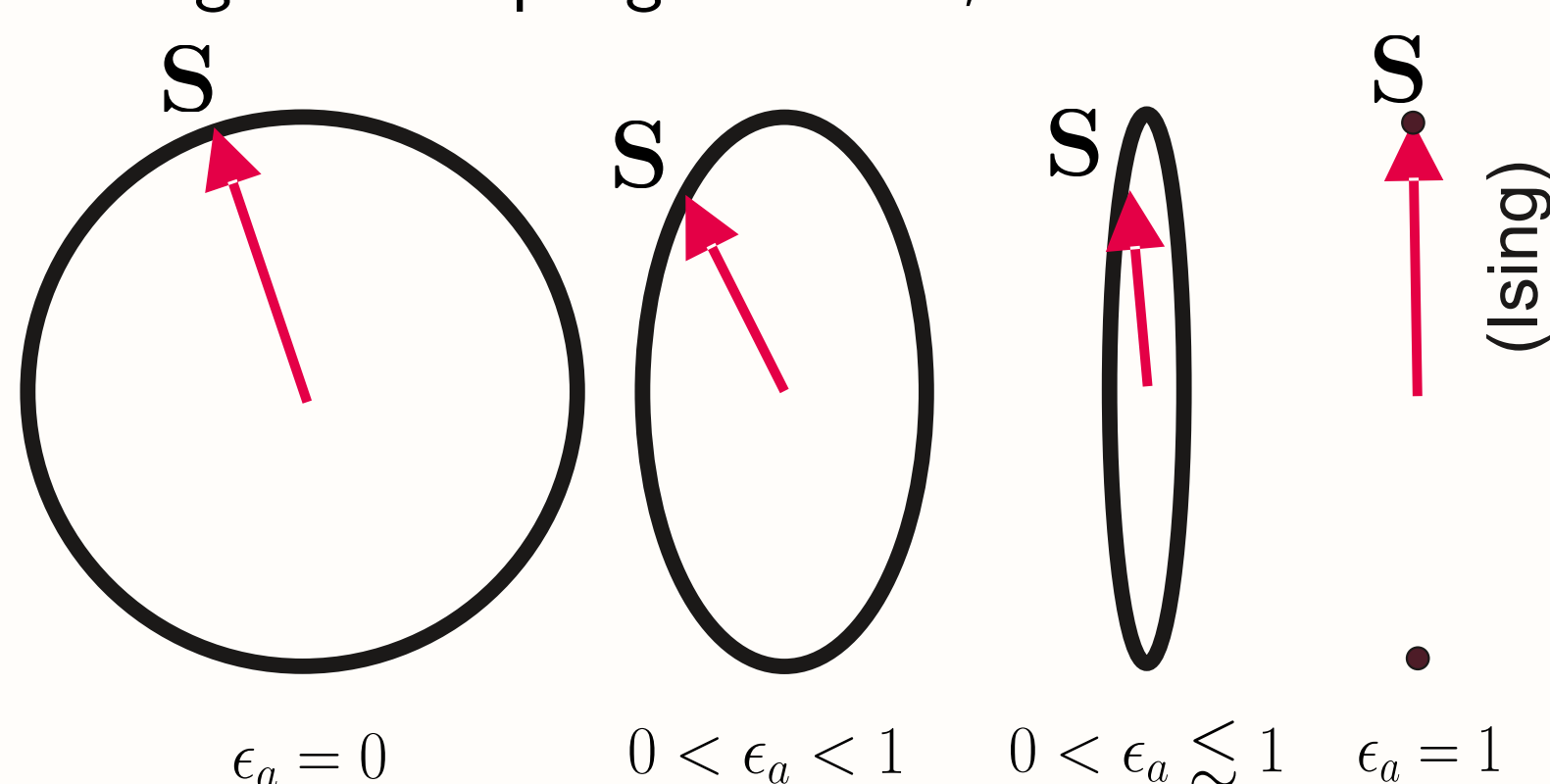


Figure 1: Schematic representation of the anisotropic XY coupling.

## Methods

- The dynamical evolution were conducted using Velocity-Verlet algorithm [4] with step size  $dt = 0.01$ ;
- The transient time :  $10^{10}$  ( $10^8$  time units) ;
- After transient the average of the heat flux is realized for  $4 \times 10^8$  time steps ( $4 \times 10^6$  time units) ;
- We set  $T_{h/l} = T(1 \pm \Delta)$  with  $\Delta = 0.125$ , where  $T$  is the average temperature ;
- The average is taken over 80 experiments ( $80 \times 4 \times 10^8$  time steps);
- The macroscopic conductivity  $\kappa$  is given by  $\kappa = \frac{J}{(T_h - T_l)/L}$  where  $J = \langle J_i \rangle$  ;

- The Lagrangian fluxes for local and anisotropic coupling are  $J_i = \frac{1}{2}(p_i + p_{i+1}) \sin(\theta_i - \theta_{i+1})$  and  $J_i = \frac{p_i + p_{i+1}}{2} \sin(\theta_i - \theta_{i+1}) + \epsilon_a \frac{p_i - p_{i+1}}{2} \sin(\theta_i + \theta_{i+1})$ , respectively.

## Results

- The present numerics at a wider range of  $T$  are proportional to stretched  $q$ -exponential whose definition  $y(x) = e_q^{-B|x|^\eta}$  with  $q \geq 1$ ,  $\eta > 0$  and  $B > 0$ . The  $q$ -Gaussian [5] form is recovered as the  $\eta = 2$  particular limit;
- Consistently with this Ansatz, we verify that, in the thermodynamic limit ( $L \gg 1$ ),  $\sigma(\epsilon_l, T) \propto \sigma(\epsilon_a, T) \propto T^{-\frac{\eta}{q-1}}$ , where  $(\eta, q) \approx (1.94, 1.65)$  thus yielding the slope  $\eta/(q-1) \approx 3.0$  (see Figs. 2 and 3).

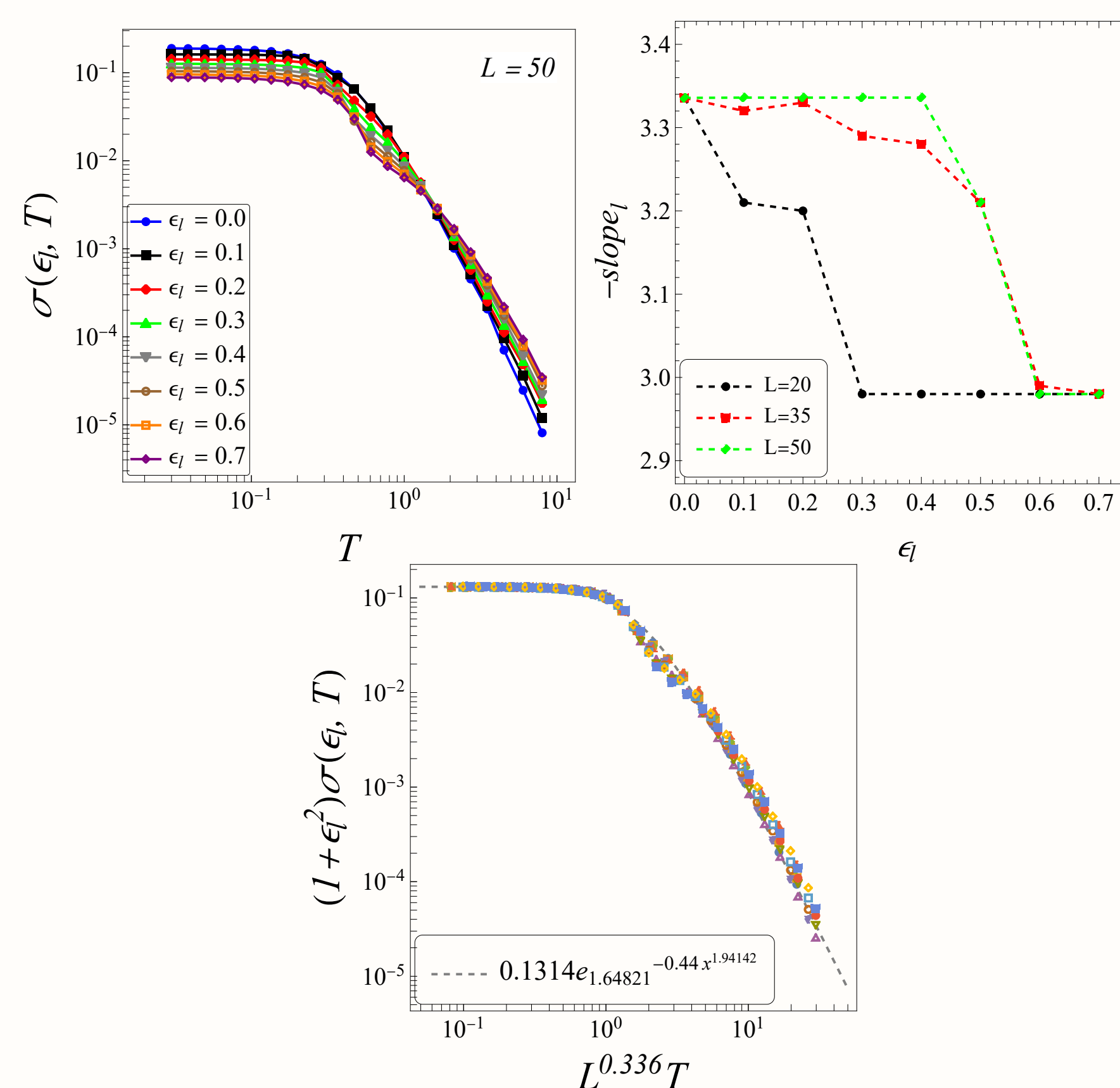


Figure 2: *Left*: Thermal conductance of the first anisotropic model as a function of temperature for one-dimensional lattice structure and the local coupling constant for  $L = 50$ . *Center*: Plot of  $-\text{slope}$  versus  $\epsilon_l$  for  $L = 20, 35, 50$ . All the curves approach the same saturation value  $\text{slope}_{\epsilon_l} \simeq -3.0$ . *Right*: Collapse with a stretched  $q$ -exponential form, from  $\epsilon_l = 0.4$  to  $\epsilon_l = 0.7$  with  $L = 20, 35, 50$ . The values of the minimum ( $T_{\min}$ ) and maximum ( $T_{\max}$ ) temperatures are 0.03 and 8.0 respectively.

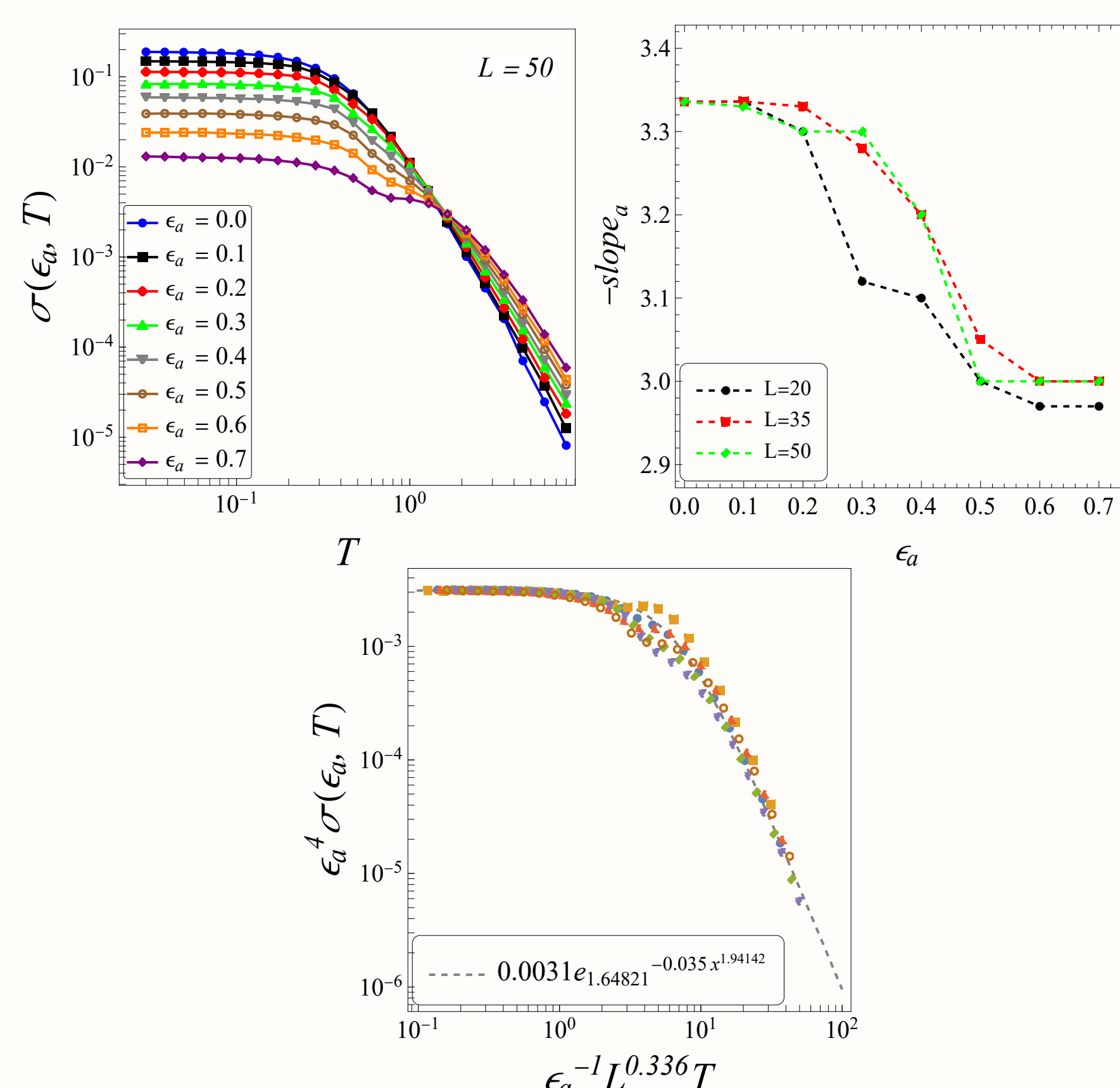


Figure 3: *Left*: Thermal conductance of the second anisotropic model as a function of temperature for one-dimensional lattice structure and the local coupling constant for  $L = 50$ . *Center*: Plot of  $-\text{slope}$  versus  $\epsilon_a$  for  $L = 20, 35, 50$ . All the curves approach the same saturation value  $\text{slope}_{\epsilon_a} \simeq -3.0$ . *Right*: Collapse with a stretched  $q$ -exponential form, for  $\epsilon_a = 0.6$  and  $\epsilon_a = 0.7$  with  $L = 20, 35, 50$ . The values of the minimum ( $T_{\min}$ ) and maximum ( $T_{\max}$ ) temperatures are 0.03 and 8.0 respectively.

## Conclusion

- Fourier's law is microscopically shown to be satisfied for the two types of anisotropic XY-models;
- A closed formula is obtained for the thermal conductivity for finite lattice sizes at arbitrary temperatures;
- In the limit of extreme anisotropy, both models approach the Ising model and its thermal conductivity  $\kappa$ , which, at high temperatures, scales like  $\kappa \sim T^{-3}$ ;
- This behavior reinforces the result obtained in various  $d$ -dimensional models, namely  $\kappa \propto L e_q^{-B(L^d T)^\eta}$  where  $e_q^z \equiv [1 + (1-q)z]^{1/(1-q)}$  ( $e_q^z = e^z$ ),  $L$  being the linear size of the  $d$ -dimensional macroscopic lattice;
- The scaling law  $\frac{\eta d}{q-1} = 1$  guarantees the validity of Fourier's law, for all dimensions.

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